

THE FOURIER TRANSFORM ON $\mathcal{F}^*(\mathbb{R}^d)$

Suppose first that $T \in \mathcal{F}^*$ is a smooth & compactly supported function, $T \in C_c^\infty(\mathbb{R}^d)$. Then

$$\begin{aligned} \langle \hat{T}, \varphi \rangle &= \int_{\mathbb{R}^d} \underbrace{\beta^d}_{\mathbb{R}^d} \int_{\mathbb{R}^d} e^{-ix \cdot t} T(x) dx \varphi(t) dt = \\ &\quad = \hat{T}(x) \boxed{\text{The integrand is absolutely summable, so we can invoke Fubini to interchange the integration order}} \\ &= \int_{\mathbb{R}^d} T(x) \underbrace{\beta^d \int e^{-ix \cdot t} \varphi(t) dt dx}_{=\hat{\varphi}(t)} = \langle T, \hat{\varphi} \rangle \end{aligned}$$

Since we've already proven that $\mathcal{F}: \mathcal{F} \rightarrow \mathcal{F}$ is continuous, we can now trivially define $\mathcal{F}: \mathcal{F}^* \rightarrow \mathcal{F}^*$ by duality.

Defn For $T \in \mathcal{F}^*(\mathbb{R}^d)$, define $\hat{T} = \mathcal{F} T$ by

$$\langle \hat{T}, \varphi \rangle = \langle T, \hat{\varphi} \rangle.$$

$$\text{Similarly, define } \overset{\vee}{T} = \mathcal{F}^* T \text{ by } \langle \overset{\vee}{T}, \varphi \rangle = \langle T, \hat{\varphi} \rangle.$$

It follows immediately that

We proved earlier that if $\varphi \in \mathcal{F}$,
then $\mathcal{F} \mathcal{F}^* \varphi = \varphi$.

$$\langle \mathcal{F}^* \mathcal{F} T, \varphi \rangle = \langle \mathcal{F} T, \mathcal{F}^* \varphi \rangle = \langle T, \overbrace{\mathcal{F} \mathcal{F}^* \varphi}^{\varphi} \rangle = \langle T, \varphi \rangle$$

So \mathcal{F}^* is the inverse of \mathcal{F} on \mathcal{F}^* . To sum up:

Propn The map $\mathcal{F}: \mathcal{F}^* \rightarrow \mathcal{F}^*$ is continuous, bijective, and ~~continuous~~ its inverse is continuous as well.

Example $T = \delta$ \leftarrow Dirac delta function

$$\langle \hat{T}, \varphi \rangle = \langle T, \hat{\varphi} \rangle = \hat{\varphi}(0) = \beta^d \int \hat{\varphi}(0) e^{-i0 \cdot x} \varphi(x) dx = \beta^d \int \varphi(x) dx = \langle \beta^d, \varphi \rangle$$

$$\text{So } \overset{\wedge}{\delta} = \frac{1}{(2\pi)^{d/2}}$$

The Fourier transform on $L^2(\mathbb{R}^d)$

For now, we switch to complex-valued functions on \mathcal{F} & \mathcal{F}^* .
 Nothing that we've done changes.

Def' Let (\cdot, \cdot) denote the ~~inner product~~ inner product, $(u, v) = \int_{\mathbb{R}^d} u(x)v(x)dx$.
 Define $\tilde{L}^2(\mathbb{R}^d)$ by taking the closure of \mathcal{F} w.r.t. the $\|\cdot\|_2$ norm.

Lemma If $\varphi, \psi \in \mathcal{F}(\mathbb{R}^d)$, then $(\varphi, \psi) = (\hat{\varphi}, \hat{\psi})$.

Proof $(\varphi, \psi) = \int \bar{\varphi} \psi = \langle \bar{\varphi}, \psi \rangle = \langle \bar{\varphi}, \hat{\psi} \rangle = \langle \frac{\vee}{\bar{\varphi}}, \hat{\psi} \rangle$ \blacksquare

$$\text{We have } \hat{\varphi}(t) = \beta^d \int e^{ixt} \bar{\varphi}(x) dx = \overline{\beta^d \int e^{-ixt} \varphi(x) dx} = \overline{\hat{\varphi}(t)} \text{ so}$$

$$(\varphi, \psi) = \langle \hat{\varphi}, \hat{\psi} \rangle = \langle \bar{\hat{\varphi}}, \hat{\psi} \rangle = (\hat{\varphi}, \hat{\psi}).$$

As a consequence of the lemma: $\|T\varphi\|_2^2 = (\hat{\varphi}, \hat{\varphi}) = (\varphi, \varphi) = \|\varphi\|_2^2$

Thus $T: \mathcal{F} \rightarrow \tilde{L}^2(\mathbb{R}^d)$ is an isometric map.

Since \mathcal{F} is dense in $L^2(\mathbb{R}^d)$, we can extend T to all of $L^2(\mathbb{R}^d)$.

Def' We define $T: L^2(\mathbb{R}^d) \rightarrow \tilde{L}^2(\mathbb{R}^d)$ as follows:

For $f \in L^2(\mathbb{R}^d)$, pick $\varphi_n \in \mathcal{F}$ s.t. $\|f - \varphi_n\|_2 \rightarrow 0$. Set $\hat{f} = \lim_{n \rightarrow \infty} \hat{\varphi}_n$

Define $T^*: \tilde{L}^2 \rightarrow L^2$ analogously, $\hat{f} = \lim_{n \rightarrow \infty} \hat{\varphi}_n$.

\hat{f} limit is in L^2

It follows immediately that $T^* = T^{-1}$ and thus that T is bijective.
 To summarize:

Thm The Fourier transform $\mathcal{F}: \mathcal{L}^2(\mathbb{R}^d) \rightarrow \mathcal{L}^2(\mathbb{R}^d)$ is

- * Bijective (one-to-one & onto)
- * Preserves the norm: $\int |u(x)|^2 = \int |\hat{u}(x)|^2$
- * Preserves the inner product: $\int \overline{u(t)} v(t) dt = \int \overline{\hat{u}(t)} \hat{v}(t) dt$
- * Is continuously invertible.

In short, it is a "unitary map" or a "Hilbert space isomorphism".

Note We defined \hat{f} as the limit $\hat{f} = \lim \hat{\phi}_n$ where $\phi_n \in \mathcal{S}$.

In practice, it is enough to pick $\phi_n \in L^2 \cap L^1$, for instance

$$\hat{f}(t) = \lim_{R \rightarrow \infty} \int_{|x| \leq R} e^{-ix \cdot t} f(x) dx = \lim_{\epsilon \rightarrow 0} \int e^{-ix \cdot t} e^{-\epsilon|x|^2} f(x) dx$$

What is the spectrum of \mathcal{F} ?

First we note that since \mathcal{F} is unitary, we must have $\sigma(\mathcal{F}) \subseteq \{\lambda \in \mathbb{C} : |\lambda|=1\}$.

Moreover, ~~indeed~~ $\mathcal{F}^2 = \mathcal{F}\mathcal{F}^* = R\mathcal{F}^*\mathcal{F} = R$, and so $\mathcal{F}^4 = R^2 = I$.

This ~~implies~~ implies that if $\lambda \in \sigma(\mathcal{F})$, then $\lambda^4 = 1$ and so $\lambda \in \{1, -1, i, -i\}$.

(Formally: $\mathcal{F}^4 = \left(\int_{\sigma(\mathcal{F})} \lambda dP(\lambda) \right)^4 = \int_{\sigma(\mathcal{F})} \lambda^4 dP(\lambda) = I \Rightarrow \lambda^4 = 1$)

In fact the four numbers $1, -1, i, -i$ are all evals of ~~mainly~~ multiplicity 1 and we have explicit formulas for the eigenvectors!

In one dimension, those are the so called Hermite functions $(\phi_n)_{n=0}^\infty$, which form an ONB for $L^2(\mathbb{R})$

$$\phi_n(x) = \alpha_n e^{x^2/2} \left(\frac{d}{dx} \right)^n e^{-x^2} = H_n(x) e^{-x^2/2} \quad \text{where } H_n \text{ are the Hermite polynomials}$$

(The H_n are constructed for instance by applying Gram-Schmidt to $\{1, x, x^2, \dots\}$ w.r.t. $\langle u, v \rangle = \int \bar{u} v e^{-x^2/2} dx$)

$$\alpha_n = \frac{1}{\sqrt{n! 2^n \sqrt{\pi}}}$$

It turns out that $\mathcal{F}\varphi_n = (-i)^\alpha \varphi_n$ and so

$$\begin{aligned} L^2(\mathbb{R}^d) &= \ker(\mathcal{F} - I) \oplus \ker(\mathcal{F} + iI) \oplus \ker(\mathcal{F} - iI) \oplus \ker(i\bar{\mathcal{F}} + I) = \\ &= \text{span}(\varphi_{4n}) \oplus \text{span}(\varphi_{4n+1}) \oplus \text{span}(\varphi_{4n+2}) \oplus \text{span}(\varphi_{4n+3}) \end{aligned}$$

With this knowledge, we could also have defined \mathcal{F} via

$$\mathcal{F}\left[\sum_{n=0}^{\infty} \alpha_n \varphi_n\right] = \sum_{n=0}^{\infty} \alpha_n (-i)^\alpha \varphi_n.$$

Sobolev Spaces

When is $\partial^\alpha \mathcal{F}$ an L^2 function?

$$\partial^\alpha \mathcal{F} \in L^2 \Leftrightarrow \int |\partial^\alpha \mathcal{F}|^2 < \infty \Leftrightarrow \int |(-it)^\alpha f(t)|^2 < \infty \Leftrightarrow (-it)^\alpha f \in L^2$$

Let s be a non-negative integer and define the Sobolev space $H^s(\mathbb{R}^d)$ by

$$\begin{aligned} H^s(\mathbb{R}^d) &= \{f : \partial^\alpha f \in L^2 \quad \forall \alpha : |\alpha| \leq s\} = \\ &= \{f : (-it)^\alpha f \in L^2 \quad \forall \alpha : |\alpha| \leq s\} = \\ &= \{f : (1+t^2)^{s/2} \hat{f} \in L^2\} = \{f : \int (1+t^2)^s |\hat{f}(t)|^2 dt < \infty\}. \end{aligned}$$

The definition is readily extended to any $s \in \mathbb{R}$, not just positive integers.

Lemma: If $f \in H^s(\mathbb{R}^d)$, for some $s > d/2$, then $f \in C_0(\mathbb{R}^d)$.

Sobolev Embedding: More generally, if $f \in H^s(\mathbb{R}^d)$ for some $s > \frac{d}{2} + k$, then $f \in C^k_c(\mathbb{R}^d)$.